Lepton flavor violation at future (lepton) colliders: induced by neutral and doubly-charged scalars

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Mini-Workshop: Theory - Physics Opportunities and Advanced Tools
IAS, HKUST

based on

P. S. B. Dev, R. N. Mohapatra & YCZ, **PRL120**(2018)221804 [1711.08430]
P. S. B. Dev, R. N. Mohapatra & YCZ, PRD**98**(2018)075028 [1803.11167]
P. S. B. Dev & YCZ, JHEP**1810**(2018)199 [1808.00943]
(see also P. S. B. Dev, M. J Ramsey-Musolf & YCZ, PRD**98**(2018)055013 [1806.08499])

see also **CEPC CDR** [1811.10545] & **CLIC Yellow Book** [1812.02093]

Outline

- Motivations of the LFV processes
- Beyond SM neutral scalar H at future lepton colliders
  - On-shell production
  - Off-shell production
  - Prospects at ILC and CEPC (CLIC in backup slides)
- Doubly-charged scalar $H^{\pm\pm}$ at future lepton colliders
  - On-shell production through the (LFV) Yukawa couplings
  - Off-shell production
  - Prospects at ILC and CEPC (CLIC in backup slides)
- Displaced LFV signals at future colliders
  - Long-lived $H_L^{\pm\pm}$ in type-II seesaw
  - DV prospects at HL-LHC, FCC-hh & ILC
  - DV from $H_R^{\pm\pm}$ in left-right models
- Conclusion

Why lepton-flavor violation (LFV) at future lepton colliders?

- Neutrino oscillations $\Rightarrow$ lepton flavor violation
  why not in the charged lepton sector???
- "Smoking-gun" signal beyond the SM;
- Clean SM background at lepton colliders, compared to the hadron colliders.

...Connection to neutrino mass generation (and other pheno)

- Beyond SM neutral scalar H from e.g. left-right model, sneutrino in RPV SUSY models;
- Doubly-charged scalar $H^{\pm\pm}$ in type-II seesaw and its extensions like left-right model;
- Might also be connected to the heavy neutrino searches, effective 4-fermion interactions, or even DM pheno at future lepton colliders.

(See the talks by R. Franceschini, J. Zupan,
M. Ramsey-Musolf, O. Fischer, M. Mitra)

Beyond SM neutral scalar H
@ future lepton colliders

Well-motivated underlying models

- RPV SUSY: sneutrinos ($\tilde{\nu}$)
  [Aulakh, Mohapatra '82; Hall,Suzuki '84; Ross, Valle '85, Barbier+ '04; Duggan, Evans, Hirschauer '13]

$$\mathcal{L}_{\rm RPV} \;=\; \frac{1}{2}\lambda_{\alpha\beta\gamma}\widehat{L}_\alpha\widehat{L}_\beta\widehat{E}^c_\gamma$$

- Left-right symmetric models: the $SU(2)_R$-breaking scalar H_3
  [Dev, Mohapatra, YCZ '16; '16; '17; Maiezza, Senjanović, Vasquez '16]

LFV couplings are generated at tree and loop level

- 2HDM: CP-even or odd (heavy) scalars from the 2nd doublet
  [Branco+ '11; Crivellin, Heeck, Stoffer '15]

LFV couplings are induced from small deviation from the lepton-specific structure.

- Mirror models: singlet scalar connecting the SM leptons to heavy mirror leptons [Hung '06, '07; Bu, Liao, Liu '08; Chang, Chang, Nugroho+ '16; Hung, Le, Tran+ '17]

LFV couplings arise from the SM-heavy lepton mixing

Effective LFV couplings

- Model-independent effective LFV couplings of H

$$\mathcal{L}_Y \;=\; h_{\alpha\beta}\bar{\ell}_{\alpha,\,L} H \ell_{\beta,\,R} \;+\; \text{H.c.}\,.$$

For simplicity, we assume $h_{\alpha\beta}$ are real, symmetric, H is CP-even, hadrophobic and the mixing with the SM Higgs h is small.

H might originate from a isospin singlet, doublet or triplet, depending on specific underlying models.

- Effective Dim-4 couplings $\neq$ Effective 4-fermion couplings like $\frac{1}{\Lambda^2}(\bar{e}e)(\bar{e}\mu)$

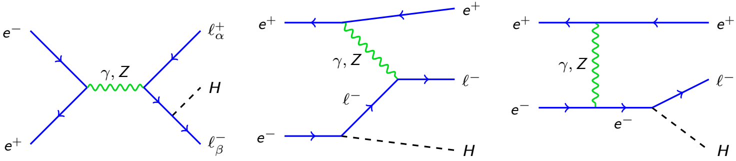
[Kabachenko, Pirogov '97; Ferreira, Guedes, Santos '06; Aranda, Flores-Tlalpa, Ramirez-Zavaleta+ '09; Murakami, Tait '14; Cho, Shimo '14]

$m_H < \sqrt{s} \Rightarrow$ on-shell production

On-shell production of H at lepton colliders

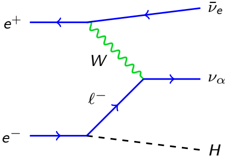
- the e^+e^- process

$$e^+e^- \to \ell_\alpha^\pm \ell_\beta^\mp + H$$



involving the charged-currents [H decaying into **visible** particles]

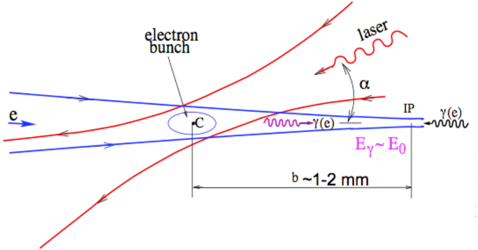
$$e^+e^- \to \nu_\alpha \bar{\nu}_e + H$$



Laser photon in future lepton colliders

- In future lepton colliders, high luminosity photon beams can be obtained by Compton backscattering of low energy, high intensity laser beam off the high energy electron beam [Ginzburg et al '83, '84].
- The effective photon luminosity distribution $\left(x = \omega/E_e \lesssim 0.83\right.$ the fraction of electron energy carried away by the scattered photon, $\left.\xi = 4\omega_0 E_e/m_e^2\right)$

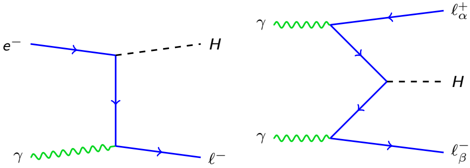
$$
\begin{aligned}
f_{\gamma/e}(x) &= \frac{1}{D(\xi)}\left[(1-x) + \frac{1}{(1-x)} - \frac{4x}{\xi(1-x)} + \frac{4x^2}{\xi^2(1-x)^2}\right], \\
\text{with} \quad D(\xi) &= \left(1 - \frac{4}{\xi} - \frac{8}{\xi^2}\right)\log(1+\xi) + \frac{1}{2} + \frac{8}{\xi} - \frac{1}{2(1+\xi)^2},
\end{aligned}
$$



On-shell production of H at lepton colliders

- involving the laser photon(s)

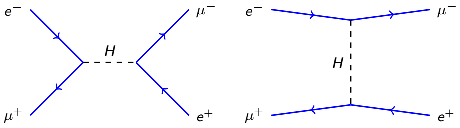
$$e^{\pm}\gamma \to \ell^{\pm} + H\,, \quad \gamma\gamma \to \ell_{\alpha}^{\pm}\ell_{\beta}^{\mp} + H$$



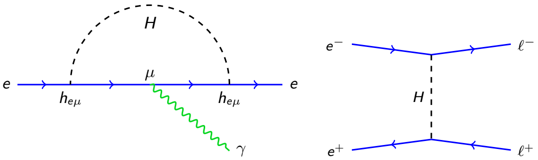
Constraints on the LFV couplings: on-shell

On-shell production amplitudes depend *linearly* on the LFV couplings

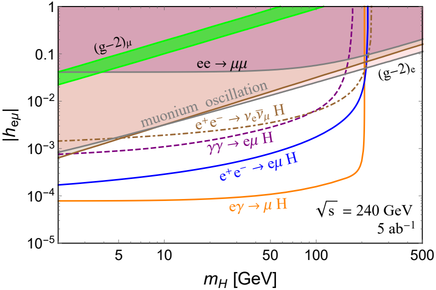
- muonium anti-muonium oscillation: $(\bar{\mu}e) \leftrightarrow (\mu\bar{e})$ $(h_{e\mu})$
  [Clark, Love '03]

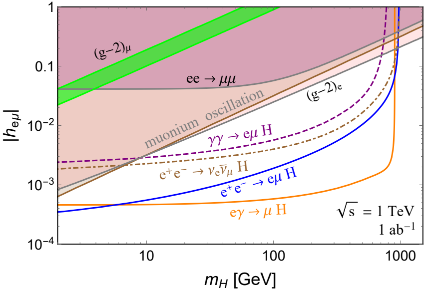


- Electron and muon $g-2$ $(h_{e\ell},\ h_{\mu\ell})$
  [Lindner, Platscher, Queiroz '16]
- Bhabha scattering, LEP $ee \to \ell\ell$ data $(h_{e\ell})$
  [OPAL '03; L3 '03; DELPHI '05]



Prospects of H: on-shell production





$\gamma\gamma$ ($e\gamma$) channel: laser photon collision.

Green bands: muon $g-2$ anomaly (excluded).

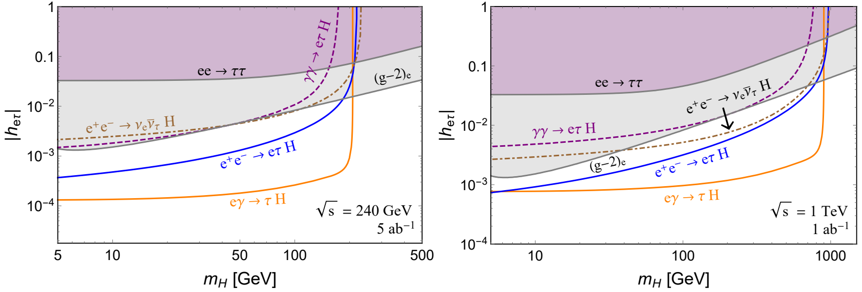
Assuming the dominant decay mode $H \to e^{\pm}\mu^{\mp}$.

Very sadly, "Japanese science committee questions the project's (ILC) multibillion-dollar price tag..."

[https://www.nature.com/articles/d41586-018-07833-9]

CLIC could do better!

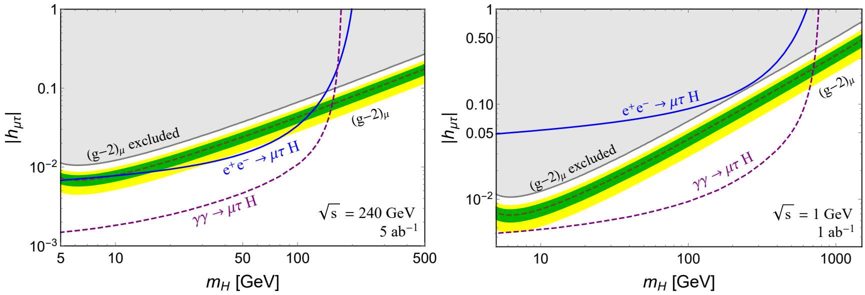
Prospects of H: on-shell production



$\gamma\gamma$ ($e\gamma$) channel: laser photon collision.

Assuming the dominant decay mode $H \to e^\pm \tau^\mp$.

Prospects of H: on-shell production



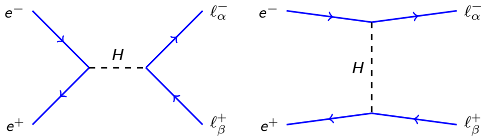
- $\gamma\gamma$ ($e\gamma$) channel: laser photon collision.
- Assuming the dominant decay mode $H \to \mu^\pm\tau^\mp$.
- The muon $g-2$ discrepancy can be directly tested at CEPC & ILC via the searches e^+e^-, $\gamma\gamma \to \mu\tau + H$.

Off-shell production of H at lepton colliders

- Off-shell production (at resonance when $m_H \simeq \sqrt{s}$)

  might also be mediated by a (light) gauge boson Z' with LFV couplings [Heeck '16]

$$e^+e^- \to \ell_\alpha^\pm \ell_\beta^\mp$$



Constraints on the LFV couplings: off-shell

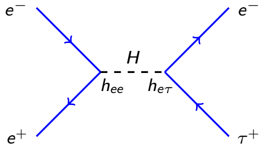
Off-shell production amplitudes depend *quadratically* on the LFV couplings

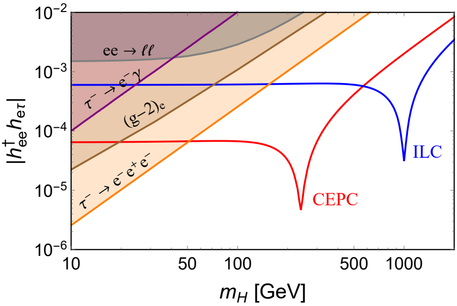
process	current data	constraints $[\mathrm{GeV}^{-2}]$
$\mu^- \to e^- e^+ e^-$	$< 10^{-12}$	$\lvert h^\dagger_{ee} h_{e\mu} \rvert / m_H^2 < 6.6 \times 10^{-11}$
$\tau^- \to e^- e^+ e^-$	$< 2.7 \times 10^{-8}$	$\lvert h^\dagger_{ee} h_{e\tau} \rvert / m_H^2 < 2.6 \times 10^{-8}$
$\tau^- \to \mu^- e^+ e^-$	$< 1.8 \times 10^{-8}$	$\lvert h^\dagger_{ee} h_{\mu\tau} \rvert / m_H^2 < 1.5 \times 10^{-8}$
$\tau^- \to \mu^+ e^- e^-$	$< 1.5 \times 10^{-8}$	$\lvert h^\dagger_{e\mu} h_{e\tau} \rvert / m_H^2 < 1.9 \times 10^{-8}$
$\tau^- \to e^- \gamma$	$< 3.3 \times 10^{-8}$	$\lvert h^\dagger_{ee} h_{e\tau} \rvert / m_H^2 < 1.0 \times 10^{-6}$
$\tau^- \to \mu^- \gamma$	$< 4.4 \times 10^{-8}$	$\lvert h^\dagger_{e\mu} h_{e\tau} \rvert / m_H^2 < 1.2 \times 10^{-6}$
$(g-2)_e$	$< 5.0 \times 10^{-13}$	$\lvert h^\dagger_{ee} h_{e\tau} \rvert / m_H^2 < 1.1 \times 10^{-7}$
		$\lvert h^\dagger_{e\mu} h_{e\tau} \rvert / m_H^2 < 1.0 \times 10^{-8}$
$ee \to ee,\ \tau\tau$	$\Lambda > 5.7\,\&\,6.3$ TeV	$\lvert h^\dagger_{ee} h_{e\tau} \rvert / m_H^2 < 1.4 \times 10^{-7}$
$ee \to \mu\mu,\ \tau\tau$	$\Lambda > 5.7\,\&\,7.9$ TeV	$\lvert h^\dagger_{e\mu} h_{e\tau} \rvert / m_H^2 < 1.3 \times 10^{-7}$

The $\mu \to 3e$ limit is so strong that the it leaves no hope to see any signal in the $ee \to e\mu$ channel at future lepton colliders.

Prospects of H: off-shell production

$$e^+e^- \to e^\pm\tau^\mp$$





- Resonance effect at $m_H \simeq \sqrt{s}$ with width $\Gamma_H = 10\,(30)$ GeV at CEPC (ILC).
- The off-shell scalar could be probed well beyond 10 TeV scale for couplings $h_{\alpha\beta}$ of order one.

Prospects of H: off-shell production

$$e^+e^- \to \mu^\pm \tau^\mp$$

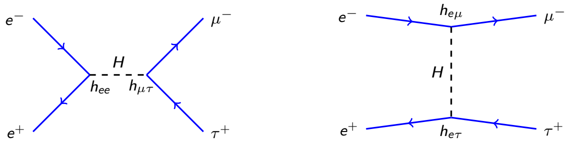
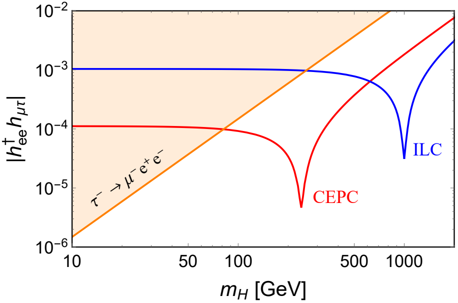
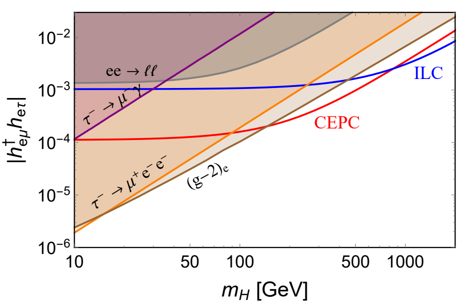


Figure: The s and t channels depend on different $h^\dagger h$ couplings.





Doubly-charged scalar $H^{\pm\pm}$

@ future lepton colliders

$H^{\pm\pm}$ at lepton (and hadron) colliders

- The (left- and right-handed) $H^{\pm\pm}$ can be pair produced from the gauge interactions to the γ/Z bosons.
- The Drell-Yan production channels can not be used to measure *directly* the (LFV) Yukawa couplings $f_{\alpha\beta}$ of $H^{\pm\pm}$ to charged leptons, unless $H^{\pm\pm}$ is long-lived.
- The current LHC same-sign dilepton limits depend largely on the branching fractions $\mathrm{BR}(H^{\pm\pm} \to \ell_\alpha^\pm \ell_\beta^\pm)$.

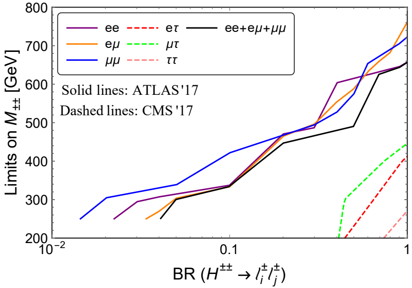


Figure: LHC dilepton limits on the right-handed $H^{\pm\pm}$.

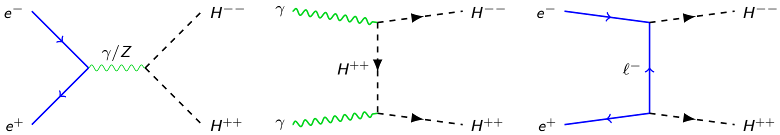
[Dev, Mohapara & YCZ, 1803.11167]

On-shell Production of $H^{\pm\pm}$ at lepton colliders through the (LFV) Yukawa couplings $f_{\alpha\beta}$

Model-independent effective couplings of (right-handed) $H^{\pm\pm}$

$$\mathcal{L}_Y = f_{\alpha\beta} H^{++} \overline{\ell^C_\alpha} \ell_\beta + \text{H.c.}$$

- Pair production through the gauge and Yukawa couplings
  [Chakrabarti+, hep-ph/9804297]

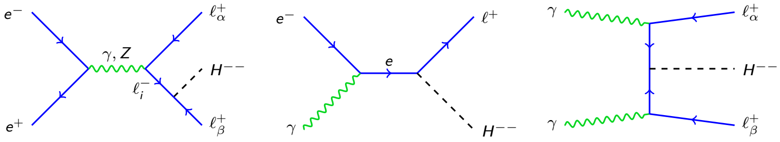


The Drell-Yan processes dominate the pair production if the Yukawa couplings $f_{e\ell}$ are very small.

On/off-shell production of $H^{\pm\pm}$ at lepton colliders

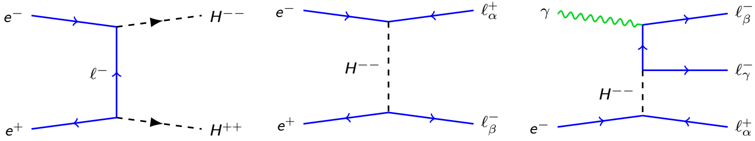
- Single production through the Yukawa couplings

  [Kuze & Sirois, hep-ex/0211048; Barenboim, Huitu, Maalampi & Raidal, hep-ph/9611362; Lusignoli & Petrarca, PLB**226**, 397; Yue & Zhao, hep-ph/0701017; Godfrey, Kalyniak, Romanenko, hep-ph/0108258; hep-ph/0207240; Rizzo, PRD**25**, 1355; Yue, Zhao & Ma, 0706.0232]

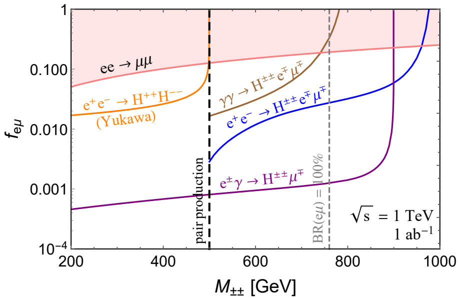


- Off-shell production

  [Godfrey, Kalyniak, Romanenko, hep-ph/0108258; hep-ph/0207240; Rizzo, PRD**25**, 1355



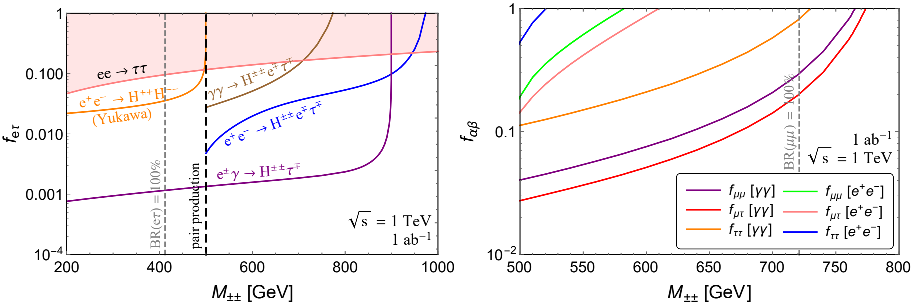
Prospects of $H^{\pm\pm}$ @ ILC 1TeV: single production



- Assuming the dominant decay mode $H^{\pm\pm} \to e^{\pm}\mu^{\pm}$.
- Below $\sqrt{s}/2 \simeq 500$ GeV, the process $e^+e^- \to H^{\pm\pm}\ell_\alpha^{\mp}\ell_\beta^{\mp}$ is dominated by the Drell-Yan pair production $e^+e^- \to H^{++}H^{--}$ with the subsequent decay $H^{\mp\mp} \to \ell_\alpha^{\mp}\ell_\beta^{\mp}$.
- The electron and muon $g-2$ limits are highly suppressed by the charge lepton masses and are not shown in the plot.

CLIC could probe higher mass ranges.

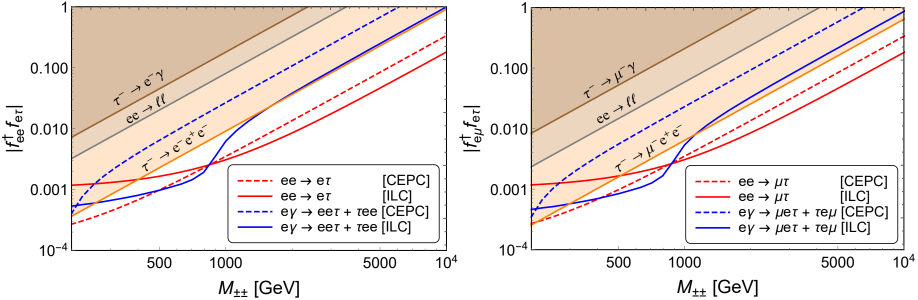
Prospects of $H^{\pm\pm}$ @ ILC 1TeV: single production



- Assuming the dominant decay mode $H^{\pm\pm} \to e^{\pm}\tau^{\pm}$ (left), $\ell_\alpha^{\pm}\ell_\beta^{\pm}$ (right).
- Below $\sqrt{s}/2 \simeq 500$ GeV, the process $e^+e^- \to H^{\pm\pm}\ell_\alpha^{\mp}\ell_\beta^{\mp}$ is dominated by the Drell-Yan pair production $e^+e^- \to H^{++}H^{--}$ with the subsequent decay $H^{\mp\mp} \to \ell_\alpha^{\mp}\ell_\beta^{\mp}$.
- The electron and muon $g-2$ limits are highly suppressed by the charge lepton masses and are not shown in the plots.

CLIC could probe higher mass ranges.

Prospects of $H^{\pm\pm}$ @ CEPC & ILC: off-shell production



- Suppressed by the three-body phase space, the sensitivities in the $e\gamma$ processes are comparatively much weaker.
- As in the neutral scalar case, the limit from $\mu \to eee$ are so stringent that it has precluded the $H^{\pm\pm}$-mediated signal $ee \to e\mu$ at CEPC & ILC.
- The effective cutoff scale $\Lambda \simeq M_{\pm\pm}/|f|$ can be probed at CEPC & ILC 1TeV up to few 10 TeV (even higher at CLIC).
- The sensitivities for more flavor combinations α, β, γ in $e^\pm\gamma \to \ell^\mp_\alpha \ell^\pm_\beta \ell^\pm_\gamma$ can be found in our paper 1803.11167.

the LFV signals might be displaced...

One example: $H_L^{\pm\pm}$ in type-II seesaw (and its right-handed partner $H_R^{\pm\pm}$ in left-right models)

Konetschny & Kummer '77; Magg & Wetterich '80; Schechter & Valle '80;
Cheng & Li '80; Mohapatra & Senjanovic '81; Lazarides, Shafi & Wetterich '81

- One of the simplest seesaw frameworks to generate the tiny neutrino masses

$$\begin{aligned} \mathcal{L} &= -(f_L)_{\alpha\beta}\,\psi^{\mathsf{T}}_{L\alpha} C i\sigma_2 \Delta_L \psi_{L\beta} + \mu H^T i\sigma_2 \Delta_L^\dagger H \;+\; \text{H.c.}, \\ \Delta_L &= \begin{pmatrix} \delta_L^+/\sqrt{2} & \delta_L^{++} = H_L^{++} \\ \delta_L^0 & -\delta_L^+/\sqrt{2} \end{pmatrix} \end{aligned}$$

- Neutrino masses are given by

$$m_\nu \;=\; \sqrt{2}\, f_L v_L \;=\; U \widehat{m}_\nu U^{\mathsf{T}} \qquad \text{(with the VEV } \langle \delta_L^0 \rangle = v_L/\sqrt{2}\text{)}$$

- The coupling matrix f_L is fixed by neutrino oscillation data, up to the unknown lightest neutrino mass m_0, the neutrino mass hierarchy, and the Dirac & Majorana CP violating phases. ⇒ VERY PREDICTIVE

Long-lived $H_L^{\pm\pm}$

- Decay through the Yukawa couplings (suppressed by m_ν^2/v_L^2)

$$\Gamma(H_L^{\pm\pm} \to \ell_\alpha^\pm \ell_\beta^\pm) = \frac{M_{H_L^{\pm\pm}}}{8\pi(1+\delta_{\alpha\beta})} \frac{|(m_\nu)_{\alpha\beta}|^2}{v_L^2},$$

- Decay through the gauge interactions
(suppressed by v_L^2 and potentially the phase-space)

$$\Gamma(H_L^{\pm\pm} \to W^\pm W^\pm) = \frac{G_F^2 v_L^2 M_{H_L^{\pm\pm}}^3}{2\pi} \sqrt{1-4x_W}(1-4x_W+12x_W^2),$$

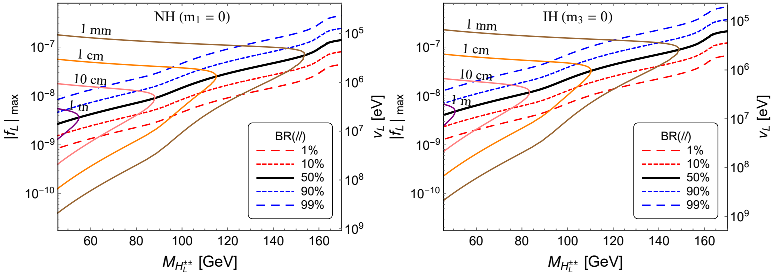
$$(\text{with } x_W \equiv m_W^2/M_{H_L^{\pm\pm}}^2)$$

Four-body decay for off-shell W-boson pairs

$$H_L^{\pm\pm} \to W^{\pm\,*} W^{\pm\,*} \to f\bar{f}' f'' \bar{f}'''$$

- Neglecting the cascade decays $H^{\pm\pm} \to H^{\pm(*)} W^{\pm(*)}$, $H^{\pm(*)} H^{\pm(*)}$ (small scalar mass splitting)

Proper lifetime of $H_L^{\pm\pm}$



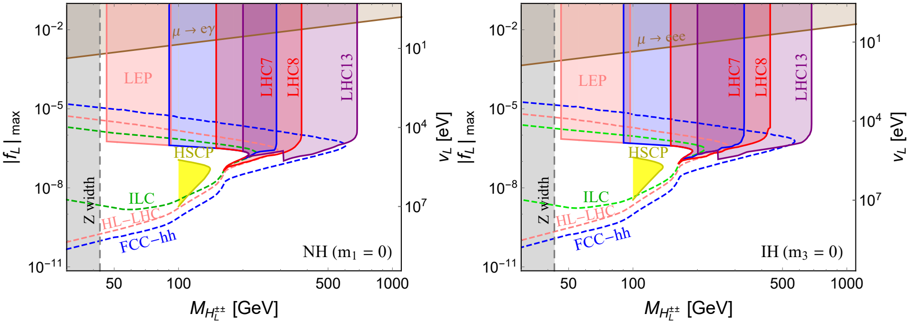
$$\Gamma_{\rm total}(H_L^{\pm\pm}) \ = \ \Gamma(H_L^{\pm\pm} \to \ell_\alpha \ell_\beta) + \Gamma(H_L^{\pm\pm} \to W^{\pm\,(*)} W^{\pm\,(*)}) \,.$$

Assuming lightest neutrino mass $m_0 = 0$.

$$v_L |f_L|_{\rm max} \ \simeq \ \begin{cases} 0.027\,{\rm eV}\,, & \text{for NH with } m_1 = 0\,, \\ 0.048\,{\rm eV}\,, & \text{for IH with } m_3 = 0\,. \end{cases}$$

Sensitivities of displaced vertices

Dev & YCZ, 1808.00943

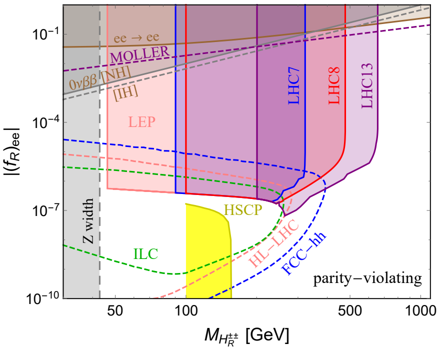


- Assuming at least 100 events for the DV sensitivities of $H_L^{\pm\pm} \to e^\pm e^\pm,\ e^\pm \mu^\pm,\ \mu^\pm \mu^\pm$.
- The low-energy high-precision LFV measurements (such as $\mu \to eee$ and $\mu \to e\gamma$), the prompt same-sign dilepton searches of $H_L^{\pm\pm}$ and the DV searches of $H_L^{\pm\pm}$ are largely complementary to each other in the type-II seesaw.

$H_R^{\pm\pm}$ in left-right symmetric model

Dev, Ramsey-Musolf & YCZ, 1806.08499

Dev & YCZ, 1808.00943



- Considering the simple scenario $H_R^{\pm\pm} \to e^\pm e^\pm$, $W_R^{\pm *} W_R^{\pm *}$.
- We do not have the LFV constraints e.g. $\mu \to e\gamma$, and MOLLER pops out...
- The low-energy high-precision LFV measurements (MOLLER and $0\nu\beta\beta$), the prompt same-sign dilepton searches of $H_R^{\pm\pm}$ and the DV searches of $H_R^{\pm\pm}$ are largely complementary to each other in the LRSM.

Conclusion

- A large variety of well-motivated models accommodate a beyond SM neutral scalar H and/or doubly-charged scalar $H^{\pm\pm}$, with LFV couplings to the SM charged leptons.
- These LFV couplings can be studied in a *model-independent* way at future lepton colliders, which strengthens the physics case for future lepton colliders.
- The neutral scalar H can be produced on-shell via $e^{\pm}\gamma \to \ell^{\pm} + H$ and $e^+e^-, \gamma\gamma \to \ell_\alpha^{\pm}\ell_\beta^{\mp} + H$ or off-shell via $e^+e^- \to \ell_\alpha^{\pm}\ell_\beta^{\mp}$.
- The doubly-charged scalar $H^{\pm\pm}$ can be (doubly & singly) on-shell and off-shell produced from the (LFV) Yukawa couplings to the charged leptons.
- It is promising that future lepton colliders could probe a broad region of mass and coupling parameters for both H and $H^{\pm\pm}$, which go well beyond the existing low-energy LFV constraints like $\tau \to eee$.
- The neutral scalar explanation of the muon $g-2$ anomaly can be directly tested at future lepton colliders in the e^+e^-, $\gamma\gamma \to \mu^{\pm}\tau^{\mp} + H$ processes.
- It might also be possible that the LFV signals are displaced, like $H^{\pm\pm} \to e^{\pm}\mu^{\pm}$ in type-II seesaw and left-right models.

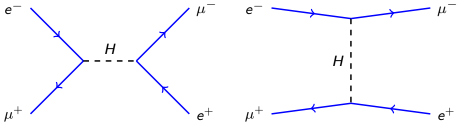
Thank you for your attention!

backup slides

Constraints on the LFV couplings $h_{\alpha\beta}$

On-shell production amplitudes depend *linearly* on the LFV couplings

- muonium anti-muonium oscillation: $(\bar{\mu}e) \leftrightarrow (\mu\bar{e})$ $(h_{e\mu})$



Oscillation probablity [Clark, Love '03]

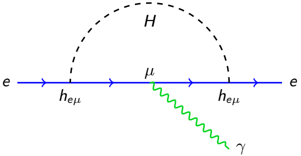
$$\mathcal{P} = \frac{2(\Delta M)^2}{\Gamma_\mu^2 + 4(\Delta M)^2}$$

with the H-induced mass splitting

$$\Delta M = \frac{2\alpha_{\rm EM}^3 h_{e\mu}^2 \mu^3}{\pi m_H^2}, \quad \mu = \frac{m_e m_\mu}{m_e + m_\mu}$$

Constraints on the LFV couplings $h_{\alpha\beta}$

- Electron and muon $g-2$ ($h_{e\ell}$, $h_{\mu\ell}$)
  [Lindner, Platscher, Queiroz '16]



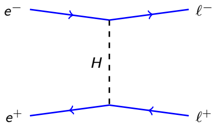
$$\Delta a_e \simeq \frac{h_{e\mu}^2 m_e m_\mu}{16\pi^2 m_H^2}\left[2\log\left(\frac{m_H^2}{m_\mu^2}\right)-3\right].$$

The value of $h_{e\mu}$ to explain $(g-2)_\mu$ discrepancy is excluded by the $(g-2)_e$ constraint.

$$\Delta a_\mu \equiv \Delta a_\mu^{\rm exp} - \Delta a_\mu^{\rm th} = (2.87 \pm 0.80) \times 10^{-9}$$

Constraints on the LFV couplings $h_{\alpha\beta}$

- Bhabha scattering, LEP $ee \to \ell\ell$ data ($h_{e\ell}$)
  [OPAL '03; L3 '03; DELPHI '05]



Effective 4-fermion interaction

$$\frac{h_{e\ell}^2}{m_H^2}(\bar{e}\ell)(\bar{e}\ell) \xrightarrow{\text{Fierz transf.}} \frac{1}{\Lambda^2}(\bar{e}\gamma_\mu e)(\bar{\ell}\gamma^\mu \ell)$$

If $m_H \lesssim \sqrt{s}$, the LEP limits on the cut-off scale Λ do not apply, and we have to consider the kinetic dependence

$$\frac{1}{m_H^2} \to \frac{1}{q^2 - m_H^2} \simeq \frac{1}{-s\cos\theta/2 - m_H^2}$$

Constraints on the LFV couplings $h_{\alpha\beta}$

Off-shell production amplitudes depend *quadratically* on the LFV couplings

- 3-body LFV decays of muon and tauon, e.g. [Sher, Yuan '91]

$$\Gamma(\tau^- \to e^+ e^- e^-) \simeq \frac{1}{\delta} \frac{|h^\dagger_{ee} h_{e\tau}|^2 m_\tau^5}{3072\pi^3 m_H^4}, \quad (\delta = 2)$$

- 2-body LFV decays of muon and tauon, e.g. [Harnik, Kopp, Zupan '12]

$$\Gamma(\tau \to e\gamma) = \frac{\alpha_{\rm EM} m_\tau^5}{64\pi^4} \left(|c_L|^2 + |c_R|^2 \right), \quad c_L = c_R \simeq \frac{h^\dagger_{ee} h_{e\tau}}{24 m_H^2}.$$

- $h_{ee,\, e\mu,\, e\tau}$ contribute to $(g-2)_e$ & LEP $ee \to \ell\ell$ data, [DELPHI '05; Hou, Wong '95]

$$|h^\dagger_{ee} h_{e\tau}| \Rightarrow ee \to e\tau$$
$$|h^\dagger_{e\mu} h_{e\tau}| \Rightarrow ee \to \mu\tau \text{ (t-channel)}$$

SM backgrounds for on-shell production of H

Main SM backgronds are particle misidentification for

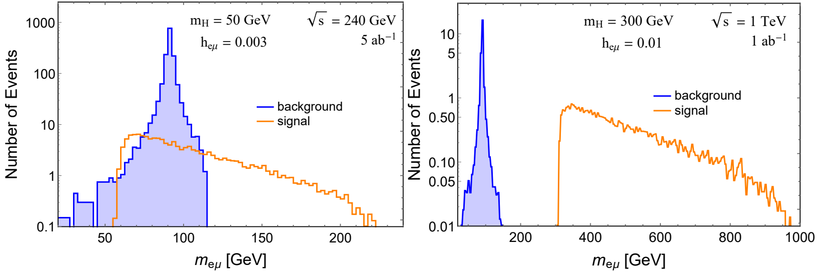
$$e^+e^- \to \ell_\alpha^+\ell_\beta^- + X\,, \quad (\alpha \neq \beta)$$

The mis-identification rate is expected to be small, of order 10^{-3}

[Milstene, Fisk, Para '06; Hammad, Khalil, Un '16; Yu, Ruan, Boudry+ '17]

Examle:

$$e^+e^- \to Zh \to (e^+e^-/\mu^+\mu^-)h \rightsquigarrow e^\pm\mu^\mp + h$$



$S/\sqrt{S+B} = 55$ $\qquad$ $S/\sqrt{S+B} = 61$

SM backgrounds for off-shell production of H

Main SM backgrounds:

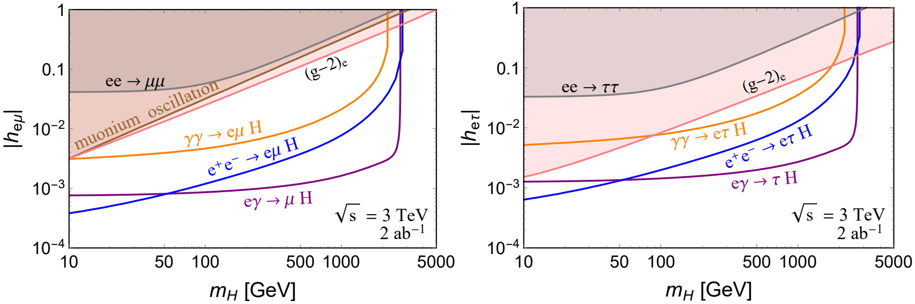
$$e^+e^- \to W^+W^- \to \ell_\alpha^+\ell_\beta^-\nu\bar{\nu}$$

The backgrounds can be well controlled by
[Kabachenko, Pirogov '97; Cho, Shimo '16; Bian, Shu, YCZ '15]

requiring that the constructed energy $E_\ell \simeq \sqrt{s}/2$,
kinetic distribution analysis of the backgrounds and signals

CLIC prospects of H and $H^{\pm\pm}$

CLIC prospects of H: on-shell production

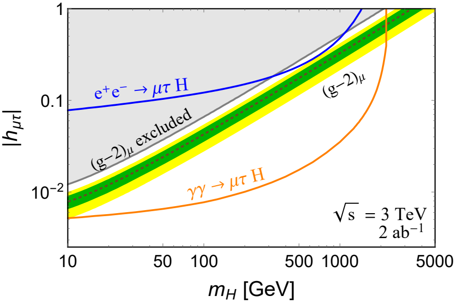


$\gamma\gamma$ ($e\gamma$) channel: laser photon collision.

Shaded regions are excluded.

Assuming the dominant decay mode $H \to e^{\pm}\mu^{\mp}$ (left), $e^{\pm}\tau^{\mp}$ (right).

CLIC prospects of H: on-shell production



$\gamma\gamma$ ($e\gamma$) channel: laser photon collision.

Shaded regions are excluded.

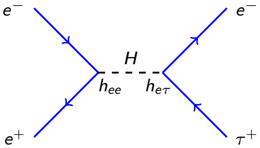
Assuming the dominant decay mode $H \to \mu^\pm \tau^\mp$.

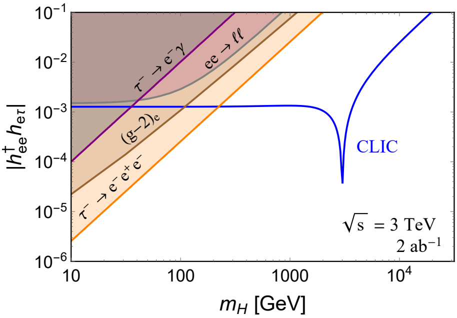
Dotted brown line: central values of muon $g-2$ anomaly, green and yellow bands: the 1σ and 2σ regions.

The muon $g-2$ discrepancy can be directly tested at CLIC via the searches of $\gamma\gamma \to \mu\tau + H$.

CLIC prospects of H: off-shell production

$$e^+e^- \to e^\pm\tau^\mp$$





Resonance effect at $m_H \simeq \sqrt{s}$ with width $\Gamma_H = 30$ GeV

The off-shell scalar could be probed well beyond 10 TeV scale (or even up to 100 TeV).

CLIC prospects of H: off-shell production

$$e^+e^- \to \mu^\pm\tau^\mp$$

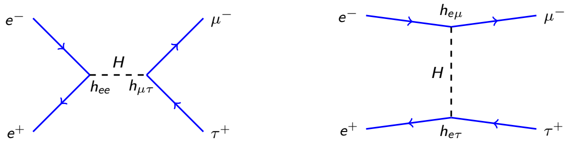
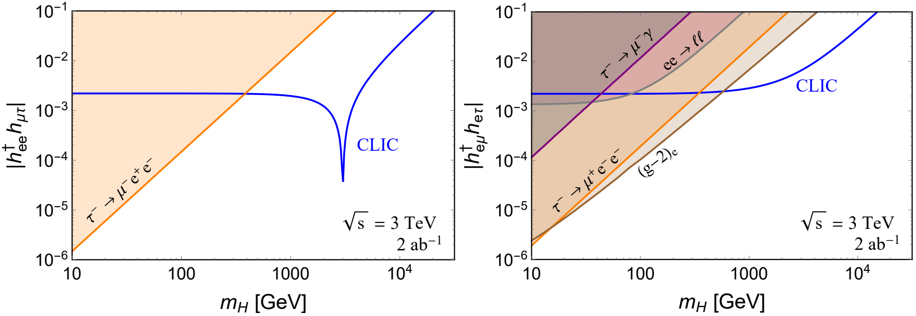
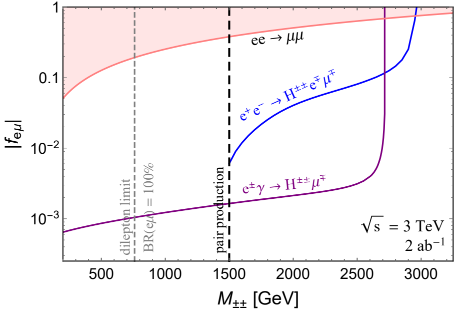


Figure: The s and t channels depend on different $h^\dagger h$ couplings.



CLIC prospects of $H^{\pm\pm}$: single production



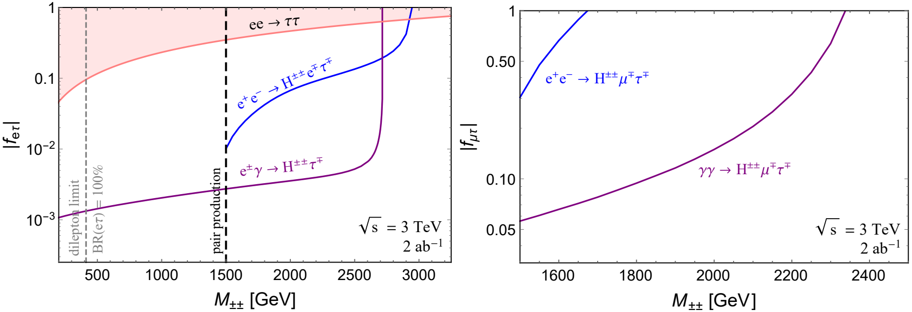
Assuming the dominant decay mode $H^{\pm\pm} \to e^\pm\mu^\pm$.

Below $\sqrt{s}/2 = 1.5$ TeV, the process $e^+e^- \to H^{\pm\pm}\ell_\alpha^\mp\ell_\beta^\mp$ is dominated by the Drell-Yan pair production $e^+e^- \to H^{++}H^{--}$ with the subsequent decay $H^{\mp\mp} \to \ell_\alpha^\mp\ell_\beta^\mp$.

The $\gamma\gamma \to H^{\pm\pm}\ell_\alpha^\mp\ell_\beta^\mp$ sensitivity is weaker than the e^+e^- process.

The electron and muon $g-2$ limits are highly suppressed by the charge lepton masses and are not shown in the plot.

CLIC prospects of $H^{\pm\pm}$: single production



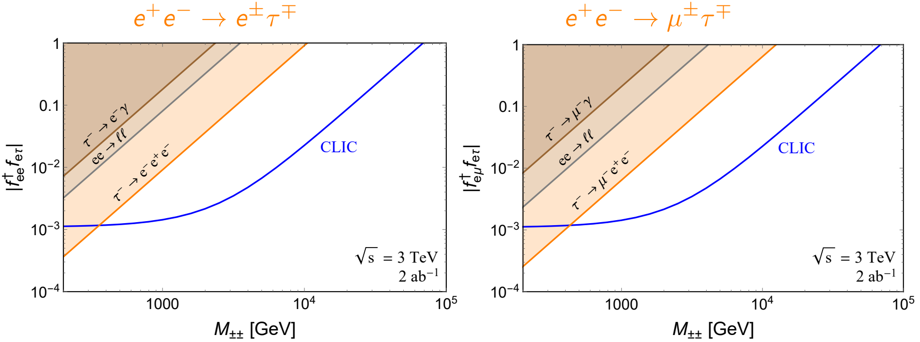
Assuming the dominant decay mode $H^{\pm\pm} \to e^\pm\tau^\pm$ (left), $\mu^\pm\tau^\pm$ (right).

Below $\sqrt{s}/2 = 1.5$ TeV, the process $e^+e^- \to H^{\pm\pm}\ell_\alpha^\mp\ell_\beta^\mp$ is dominated by the Drell-Yan pair production $e^+e^- \to H^{++}H^{--}$ with the subsequent decay $H^{\mp\mp} \to \ell_\alpha^\mp\ell_\beta^\mp$.

The $\gamma\gamma \to H^{\pm\pm}\ell_\alpha^\mp\ell_\beta^\mp$ sensitivity is weaker than the e^+e^- process.

The electron and muon $g-2$ limits are highly suppressed by the charge lepton masses and are not shown in the plots.

CLIC prospects of $H^{\pm\pm}$: off-shell production



Suppressed by the three-body phase space, the sensitivities in the $e\gamma$ processes are comparatively much weaker.

As in the neutral scalar case, the limit from $\mu \to eee$ are so stringent that it has precluded the $H^{\pm\pm}$-mediated signal $ee \to e\mu$ at CLIC.

The effective cutoff scale $\Lambda \simeq M_{\pm\pm}/|f|$ can be probed at CLIC up to few 10 TeV.